

# Elastic scattering using the optical model within NRV Knowledge Base

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To understand:

- 1 Mathematical formulation of low-energy quantum elastic scattering
- 2 Use NRV Knowledge Base in describing elastic scattering
- 3 How to describe elastic scattering in NRV:
  - ${}^4\text{He} + {}^{58}\text{Ni} \rightarrow {}^4\text{He} + {}^{58}\text{Ni}$
  - ${}^4\text{He} + {}^{209}\text{Bi} \rightarrow {}^4\text{He} + {}^{209}\text{Bi}$at different energies

## Introduction: Nuclear Reactions Video (NRV) Knowledge Base

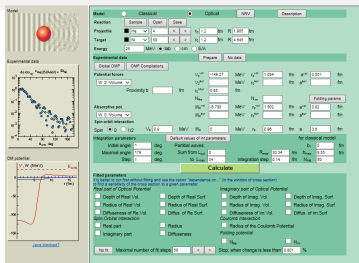
- 1 NRV: Developed at Flerov Laboratory of Nuclear Reactions, Joint Institute for Nuclear Research
- 2 Allows one to simulate many channels of nuclear reactions at low energies: Elastic scattering, inelastic scattering, transfer reactions, fusion, etc.



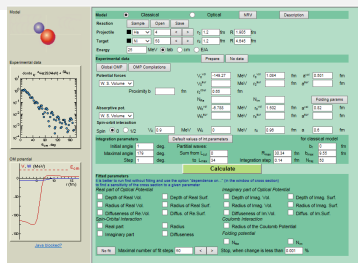
Figure: NRV homepage. URL: <http://nrv.jinr.ru>

- 3 In this project: elastic scattering of nuclei within the NRV project.

# NRV Screenshots, Elastic scattering: Quantum OM and Classical Model



(a) Input page, Quantum OM



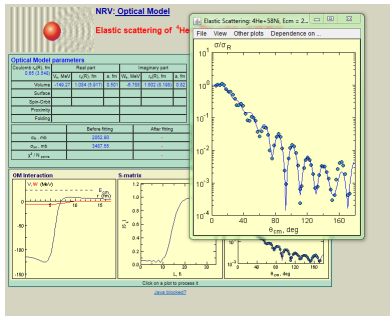
(b) Input page, Classical.

## Input information

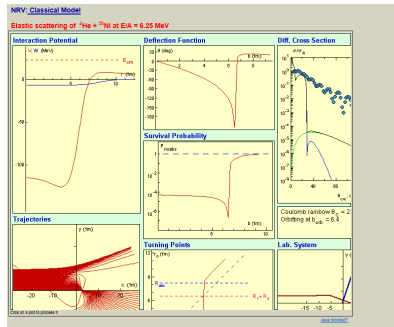
Figure

- 1 Reaction:
  - The projectile and target selected from the drop-down lists
  - Mass number of the projectile and target
  - Energy (lab., centre-of-mass or energy per nucleon of projectile)
- 2 Experimental data, if available. May be copied into NRV from a text file
- 3 Potential: A real and imaginary part of the optical potential; spin-orbit potential, if the projectile has a spin. Automatic values available for potential parameters.
- 4 Integration parameters

# NRV results page



(a) Results page, OM:  $\sigma$ ,  $d\sigma/d\Omega$ ,  $S_{\ell}(k)$ , partial wavefunctions, etc



(b) Results page, Classical:  $\sigma$ ,  $d\sigma/d\Omega$ , trajectories field, deflection, turning point, survival prob., etc

Figure

## Quantum elastic scattering theory

- 1 Elastic scattering:  $A + B \rightarrow A + B$
- 2 Time-independent formalism: stationary states
- 3 Time-Independent Schrödinger Equation:  $H\psi^+ = E\psi^+$ ,  $E > 0$
- 4 Boundary conditions
- 5 Spherically symmetric domain, total wavefunction  $\psi(r, \theta, \phi)$  separable:  
$$\psi(r, \theta, \phi) = \underbrace{(\psi_\ell(r)/r)}_{\text{radial}} \underbrace{Y_\ell^m(\theta, \phi)}_{\text{spherical}}$$
- 6 Azimuthal symmetry,  $m = 0$ ,  
 $Y_\ell^0(\theta, \phi) = \sqrt{(2\ell + 1)/4\pi} P_\ell(\cos\theta)$
- 7  $\psi_\ell(r)$  satisfies the single-channel Radial Schrödinger Equation:

$$-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} \psi_\ell + \left[ V(r) + \frac{\ell(\ell + 1)\hbar^2}{2\mu r^2} \right] \psi_\ell = E_{cm} \psi_\ell \quad (1)$$

$V(r)$  is the potential.

- 1 At  $r = 0$

$$\psi_{k\ell}^+(r) = 0$$

- 2 Asymptotic region (Outgoing waves boundary conditions)

$$\psi_{k\ell}^+(r) \rightarrow \frac{i}{2}[H_{\ell}^-(kr; \eta) - S_{\ell}(k)H_{\ell}^+(kr; \eta)], \quad r \rightarrow \infty \quad (2)$$

where  $S_{\ell}(k) = e^{2i(\sigma_{\ell}(k) + \delta_{\ell}^N(k))}$ ,  $\eta$  the Sommerfeld parameter,  $H_{\ell}^-(kr; \eta)$  and  $H_{\ell}^+(kr; \eta)$  are the incoming and outgoing waves, respectively.

- 3  $H_{\ell}^{\pm}(kr; \eta) = G_{\ell}(kr; \eta) \pm iF_{\ell}(kr; \eta)$ ,  
where  $G_{\ell}(kr; \eta)$  and  $F_{\ell}(kr; \eta)$  are the regular and irregular Coulomb functions, resp.; two linearly independent solutions of pure Coulomb scattering.
- 4 Numerov Method is used in NRV to solve the Radial Schrödinger Equation

- 1 Partial wave analysis (Faxen-Holzmark formalism):

$$f(\theta) = \frac{1}{2ik} \sum_{\ell=0}^{\infty} (S_{\ell}(k) - 1)[(2\ell + 1)P_{\ell}(\cos \theta)] \quad (3)$$

$S_{\ell}(k)$  contains information on nature of the interaction: short range, long range

- 2  $f(\theta) = f_C(\theta) + f_N(\theta)$
- 3  $f(\theta)$  establishes the link between scattering theory and experiments.

Differential cross-section: (4)

$$\frac{d\sigma}{d\Omega} = |f(\theta, \phi)|^2 \quad (5)$$

Integrated cross-section: (6)

$$\sigma = 4\pi \sum_{\ell=0}^{\infty} (2\ell + 1) |f_{\ell}|^2 \quad (7)$$

where  $f_{\ell}$  are the partial scattering amplitudes.



- ① 
$$V(r) = \underbrace{V_C(r)}_{\text{Coulomb}} + \underbrace{V_N(r) + iW_N(r)}_{\text{Nuclear}} + \underbrace{V_{SO}(r) + iW_{SO}(r)(\vec{l} \cdot \vec{s})}_{\text{Spin-orbit}}$$
- ② Coulomb potential for uniform charge distribution:

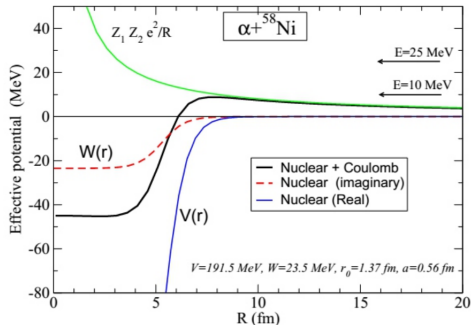
$$V_C(r) = \begin{cases} Z_P Z_T e^2 \cdot \frac{1}{2R_C} \cdot (3 - \frac{r^2}{R_C^2}), & \text{for } r < R_C \\ Z_P Z_T e^2 \cdot \frac{1}{r}, & \text{for } r > R_C \end{cases} \quad (8)$$

- ③  $V_N(r)$  and  $W_N(r)$  may be approximated as follows:
- Woods-Saxon Volume (WSV) potential
  - Woods-Saxon Surface (WSS) potential
  - Superposition of WSV and WSS
  - Folding Potential

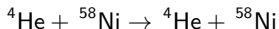
## Woods-Saxon Volume (WSV) potential

- In this project WSV potentials are used for both  $V_N(r)$  and  $W_N(r)$ :

$$V_N(r) = -\frac{V_o}{1 + \exp[(r - r_v)/a_v]}, \quad W_N(r) = -\frac{W_o}{1 + \exp[(r - r_w)/a_w]}$$

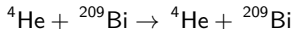


- Two simulations: alpha particle + heavy-ion  
 ${}^4\text{He}$  has zero spin, no spin-orbit interaction,  $V_{SO} = W_{SO} = 0$



Coulomb barrier = 11.50 MeV(*c.m.*)

$E_{cm} = 7.6, 8.5, \text{ and } 23.4 \text{ MeV}$



Coulomb barrier = 26.00 MeV(*c.m.*)

$E_{cm} = 24.3, 34.0, \text{ and } 68.2 \text{ MeV}$

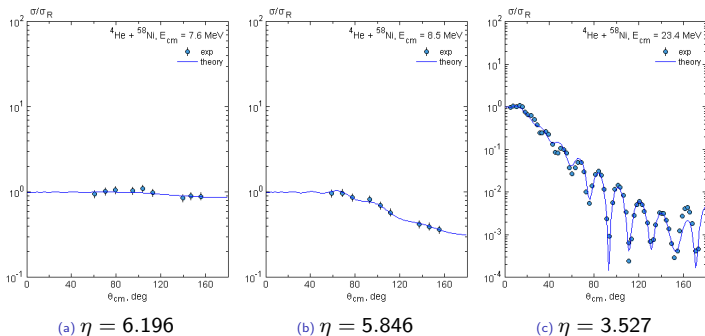


Figure: Ratio of total cross-section to Rutherford cross-section. Rutherford scattering at low energies; Fresnel behaviour around barrier, Fraunhofer oscillations for high energies beyond Coulomb barrier.

$$\eta = \frac{Z_P Z_T e^2}{\hbar} \left( \frac{\mu}{2E} \right)^{1/2}; \text{ Nickel-58, } Z=28$$

$E_{cm}/\text{MeV}$	Real (Woods-Saxon Vol.)			Imaginary (Woods-Saxon Vol.)		
	$V_o/\text{MeV}$	$r_v/\text{fm}$	$a_v/\text{fm}$	$W_o$	$r_w/\text{fm}$	$a_w/\text{fm}$
7.6	-155.106	1.068	0.53	-14.886	0.479	0.820
8.5	-154.134	1.027	0.676	-16.866	0.306	1.649
23.4	-149.270	1.084	0.501	-6.788	1.502	0.820

Table: Wood-Saxon Vol. parameters for the nucleus-nucleus potential

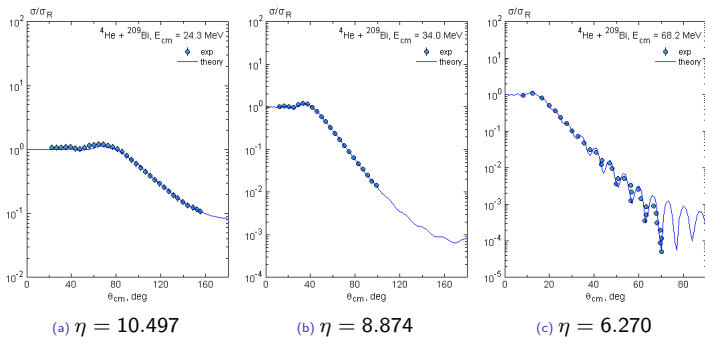


Figure: Ratio of total cross-section to Rutherford cross-section. Fresnel behaviour around barrier, Fraunhofer oscillations for high energies beyond Coulomb barrier.

$$\eta = \frac{Z_P Z_T e^2}{\hbar} \left( \frac{\mu}{2E} \right)^{1/2}; \text{ Bismuth-209, } Z = 83$$

$E_{cm}/\text{MeV}$	Real (Woods-Saxon Vol.)			Imaginary (Woods-Saxon Vol.)		
	$V_o/\text{MeV}$	$r_v/\text{fm}$	$a_v/\text{fm}$	$W_o$	$r_w/\text{fm}$	$a_w/\text{fm}$
24.3	-162.9	0.966	0.656	-11.478	0.979	0.472
34.0	-163.0	1.035	0.593	-17.128	1.134	0.400
68.2	-163.5	1.014	0.656	-17.800	1.153	0.600

Table: Wood-Saxon Vol. parameters for the nucleus-nucleus potential

- 1 Elastic scattering studied within the Optical Model
- 2 NRV OM code used to study
$${}^4\text{He} + {}^{58}\text{Ni}$$
$${}^4\text{He} + {}^{209}\text{Bi}$$
at different energies
- 3 We have obtained a good agreement between the calculations and experimental data